ANDREW R. WHITEHILL<sup>1,\*</sup>, CHANGJIAN XIE<sup>2</sup>, XIXI HU<sup>2</sup>, DANQIAN XIE<sup>2</sup>, HUA GUO<sup>3</sup> AND SHUHEI ONO<sup>1</sup>

<sup>1</sup>Department of Earth, Atmospheric, and Planetary Sciences, Massachusetts Institute of Technology, Cambridge, MA

<sup>2</sup>Institute of Theoretical and Computational Chemistry, School of Chemistry and Chemical Engineering, Nanjing University, Nanjing, China

<sup>3</sup>Department of Chemistry and Chemical Biology, University of New Mexico, Alburquerque, NM, USA \*Correspondance to: arwhite@mit.edu

Sulfur dioxide is known to produce different patterns of sulfur mass-independent isotope fractionation (S-MIF) in two different absorption bands. SO<sub>2</sub> photolysis at wavelengths between 180 and 220 nm produces significant S-MIF only at high optical densities (>10<sup>-16</sup> molecules/cm<sup>2</sup>), suggesting a dominant contribution from optical shielding effects[1]. Both mass-dependent and mass-independent isotope effects become larger at higher SO<sub>2</sub> column densities (up to optical saturation) and at lower temperatures. Mass-dependent isotope effects are consistent with cross-section calculations. SO<sub>2</sub> photolysis occurs due to coupling of a bound state (<sup>1</sup>B<sub>2</sub>) with the dissociative continuum of the ground state (<sup>1</sup>A<sub>1</sub>) [2]. The high density of states in the <sup>1</sup>A<sub>1</sub> state makes it unlikely that there is a considerable difference in photolysis quantum yields between isotopologues.

In contrast, photoexcitation of SO<sub>2</sub> in the 250 to 350 nm region produces S-MIF that is not due to self-shielding. Its pattern differs considerably from that predicted by cross-sections. Photoexcitation produces larger S-MIF signatures at lower SO<sub>2</sub> pressures and produces very large S-MIF signatures even under optically thin conditions. The isotope effects in this absorption region are due to an isotopologue-specific dependence on intersystem crossing rates between the excited singlet (coupled  ${}^{1}B_{1}{}^{1}A_{2}$ ) states and reactive triplet ( ${}^{3}B_{1}$ ) states [3], both of which have a low density of states in the crossing region.

Understanding the mechanisms for the production of massindependent fractionation allows us to predict other molecules that might display similar effects (e.g.  $CS_2$  and SO) and explain the mechanism responsible for the production of S-MIF signatures in the modern atmosphere [4].

[1] Ono et al (2013) J. Geophys. Res. Atmos. 118, 2444 [2]
Katagiri et al (1997) J. Molec. Struct. 413-414, 589 [3]
Whitehill et al (2013) PNAS 110, 17697 [4] Savarino et al (2003) Geophys. Res. Lett. 30, 2131